



## Communication

Capture and separation of CO<sub>2</sub> on BC<sub>3</sub> nanosheets: A DFT study

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## ABSTRACT

In order to reduce the greenhouse effect caused by the rapid increase of CO<sub>2</sub> concentration in the atmosphere, it is necessary to develop more efficient, controllable, and highly sensitive adsorbing materials. In this study, the adsorption behavior of CO<sub>2</sub> on BC<sub>3</sub> nanosheets under an external electric field was explored based on density functional theory (DFT). It was found that CO<sub>2</sub> experienced a transition from physisorption to chemisorption in the electric field range of 0.0060–0.0065 a.u.. In addition, the adsorption/desorption of CO<sub>2</sub> is reversible and can be precisely controlled by switching on/off at the electric field of 0.0065 a.u.. The selective adsorption of CO<sub>2</sub>/H<sub>2</sub>/CH<sub>4</sub> by BC<sub>3</sub> can also be used to realize gas separation and purification under different electric fields. This study highlighted the potential application of BC<sub>3</sub> nanosheets as a high-performance, controllable material for CO<sub>2</sub> capture, regeneration, and separation in an electric field.

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CO<sub>2</sub> capture and storage has attracted extensive research due to the greenhouse effect caused by its concentration rapid increase, which caused environmental deterioration and global warming. It is necessary to consider the nature of the adsorbent materials containing the high adsorption capacity, sensitive selectivity [1], low energy consumption during adsorption/desorption process, more reusable times [2,3] and structural stability [4]. At present, the popular CO<sub>2</sub> adsorbents are the metal-organic frameworks (MOFs) [5,6], covalent organic frameworks (COFs) [7], 2D solid materials including graphene [8], graphene-like materials [9], porous carbon materials and polymers [10–13].

In the past few years, many researches have been done on the use of external electric fields to promote CO<sub>2</sub> adsorption and achieve CO<sub>2</sub> separation. Guo *et al.* found that the hexagonal BN nanosheets have high-efficiency adsorption and large capacity for CO<sub>2</sub> under an applied electric field, which is effectively separate CO<sub>2</sub> from the CO<sub>2</sub>/H<sub>2</sub>/N<sub>2</sub>/CH<sub>4</sub>/CO/H<sub>2</sub>O mixture [14]. The work of Adnan *et al.* shows that under an external electric field of 0.013 a.u.,

CO<sub>2</sub> achieves the transition from physisorption to chemisorption, while the desorption of CO<sub>2</sub> is a spontaneous exothermic process with 0.06 eV after the removal of electric field [13]. Qiao *et al.* reports that CO<sub>2</sub> capture/regeneration on MoS<sub>2</sub> monolayers can be controlled by external electric fields, and the adsorbed CO<sub>2</sub> can be released without any energy barrier once the electric field is turned off [15].

Recently, BC<sub>3</sub> monolayer with high crystalline quality has been successfully created by carbon-substituted technique in a boron hexagonal comb [16]. BC<sub>3</sub> has excellent mechanical properties [17,18] and thermal stability [19]. It has excellent performance in catalytic reactions [20], and the potential application for gas sensors [21,22].

The first principle is the ab initio calculation, which does not need any parameters. It only needs to start from the most basic physical principle to directly solve the Schrodinger equation with some approximated algorithm, according to the interactions between nucleus and electron and its basic motion rules. The first principle is widely used to calculate the catalysis process of single atom supported or bimetallic catalysts [23–25], chemical reactions caused by external electrical fields [26,27], surface adsorption and activity [28], organics degradation simulation [29], analysis of photocatalytic mechanism [30,31], optical properties prediction [32] and so on [33–45].

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In this research, we systematically explored whether BC<sub>3</sub> has the potential to control capture and release of CO<sub>2</sub> under an external electric field. First, the adsorption of CO<sub>2</sub>/H<sub>2</sub>/CH<sub>4</sub> on BC<sub>3</sub> surface was studied, then the transformation of CO<sub>2</sub> adsorption behavior under gradient electric field and the thermodynamic process of CO<sub>2</sub> adsorption/desorption on the surface of BC<sub>3</sub> were further explored. Finally, the application of different electric field values to separate CO<sub>2</sub>/CH<sub>4</sub>/H<sub>2</sub> mixed gas was evaluated.

All of the spin-polarized DFT calculations were performed using the Dmol<sup>3</sup> [46] software. Generalized gradient approximation (GGA) formulated by Perdew, Burke, and Ernzerhof (PBE) are carried out to describe the exchange-correlation energy functional [47], and meanwhile we used all-electron core treatment for the electron ion-core interactions [48], and a double numerical basis set containing d polarization function, such as double numeric with polarization (DNP) to increase accuracy. The spin polarization is turned on all the way, and Grimme method was applied to describe the van der Waals interactions [49]. The charge distributions and the charge transfer between BC<sub>3</sub> monolayer and molecules were analyzed by using Hirshfeld method. For geometry optimization and electronic self-consistency, the convergence tolerance is set to  $1 \times 10^{-6}$  a.u. The convergence threshold parameters of geometric optimization are specified as  $1 \times 10^{-5}$  Ha,  $2 \times 10^{-3}$  Ha/Å and  $5 \times 10^{-3}$  Å for energy, maximum force and displacement, respectively. In the process of structural optimization, the coordinates of the internal atoms were relaxed with the fixed lattice parameters.

The Brillouin k-point grid was sampled by the  $6 \times 6 \times 1$  Monkhorst pack grid [50]. The BC<sub>3</sub> base model used  $2 \times 2$  supercell containing 32 atoms with a 20 Å vacuum layer to weaken interaction between layers. The optimized lattice constant of the BC<sub>3</sub> structure is 10.35 Å, which is well consistent with the theoretical result of 10.33 Å [51]. Then the adsorption energy was calculated by this formula (Eq. 1):

$$E_{\text{ads}} = E_{\text{All}} - E_{\text{sub}} - E_{\text{gas}} \quad (1)$$

Where  $E_{\text{All}}$  is the total energy of the absorbed combination,  $E_{\text{sub}}$  and  $E_{\text{gas}}$  represent the energy of separate BC<sub>3</sub> substrate and gas molecules, respectively. It is worth noting that the E-field perpendiculars to the BC<sub>3</sub> surface and points downward.

As shown in Fig. S1a (Supporting information), BC<sub>3</sub> nanosheets has a planar hexagonal honeycomb structure, each structural unit includes four C atoms and two B atoms. Each B atom was combined with three C atoms via  $sp^2$  hybridization. The lengths of B–C and C–C are 1.565 and 1.422 Å, respectively, which was well consistent with the early literature [52]. As depicted in the Fig. S1a, six adsorption sites are considered for molecular adsorption, including the top sites of B (T1) and C (T2) atoms, the bridge sites of C–C (B1) and B–C bond (B2), and the central position of C<sub>4</sub>B<sub>2</sub> six-membered ring (H1) and the C six-membered ring (H2).

The structural stability of BC<sub>3</sub> nanosheets under the application of an E-field can be roughly evaluated by the polymerization energy and molecular dynamics simulation. The formula for polymerization energy was defined as Eq. 2:

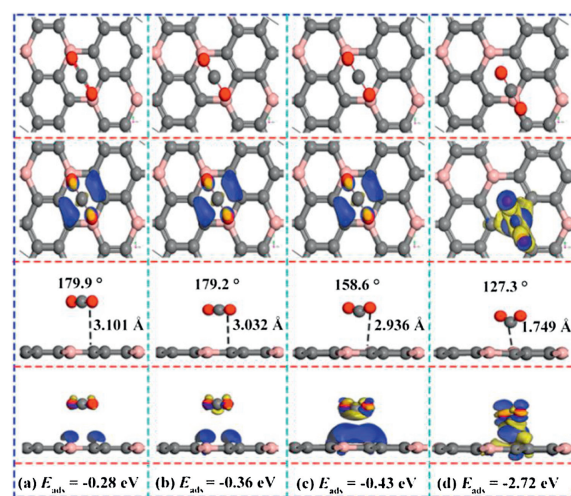
$$E_{\text{coh}} = (nE_{\text{B}} + mE_{\text{C}} - E_{\text{BC}_3}) / (m + n) \quad (2)$$

Where  $E_{\text{B}}$ ,  $E_{\text{C}}$ ,  $E_{\text{BC}_3}$  represents the energy of B and C atoms and BC<sub>3</sub> nanosheets, respectively,  $n$  and  $m$  are the number of C and B atoms [53]. As shown in Fig. S1b (Supporting information), as the E-field strength increased from 0 to 0.008 a.u., the average binding energy decreased from 7.52 eV to 7.40 eV. Compared with that of no E-field, the average binding energy between BC<sub>3</sub> atoms changes little, indicating that the structure is stable. For the optimized structure under the electric field the microcanonical ensemble (NVE) ensemble was adopted, at 300 K, molecular dynamics simulation

was done. As shown in Fig. S1c (Supporting information), the energy change of the structure is very small in the 2 ps interval of 1 fs, which also confirmed the stability of the BC<sub>3</sub> structure under the applied E-field.

We studied three molecules (CO<sub>2</sub>, CH<sub>4</sub>, H<sub>2</sub>) adsorption on BC<sub>3</sub> substrate, the most stable adsorption structure and relative adsorption parameters were depicted in Fig. S2 and Table S1 (Supporting information), respectively. As shown in Fig. S2, CO<sub>2</sub> tends to adsorb on H1 site of C<sub>4</sub>B<sub>2</sub> six-membered ring, however, CH<sub>4</sub> and H<sub>2</sub> tend to adsorb on T2 and T1 sites, respectively. No chemical bond is formed between the molecules and the substrate, where the CO<sub>2</sub>, CH<sub>4</sub>, H<sub>2</sub> are suspended above the surface. The minimum distance ( $D_m$ ) between molecules and BC<sub>3</sub> substrate are 3.271, 3.609 and 3.158 Å for CH<sub>4</sub>, CO<sub>2</sub> and H<sub>2</sub>, respectively, and the corresponding adsorption energy is very small with the value of  $-0.11$ ,  $-0.09$  and  $-0.07$  eV, respectively. From the analysis of Hirshfeld-type charge transfer, the amount of charge transfer between gas molecules and BC<sub>3</sub> is also very small (0.004 e,  $-0.025$  e and 0.012 e).

As we all know, the strong interactions of the adsorbates combined with surface contribute to the evident results of gas detection selectivity [54] and sensitivity. In order to enhance the CO<sub>2</sub> sensitivity, we apply a forward E-field perpendicular to the Z axis [55] on BC<sub>3</sub> surface [52]. The corresponding adsorption configurations were presented in Fig. 1. According to the difference charge density, the adsorption effect gradually increases along with the increasing of the electric field strength. As shown in Fig. 1c, the angle of  $\angle$ O–C–O decreased sharply from 180° to 158.6° along with the electric field increased from 0 to 0.006 a.u., and the adsorption distance reduced from 3.10 Å to 2.94 Å. The adsorption energy adds up to  $-0.43$  eV, and the charge transfer amount increases to 0.2 e, the charge density expands a lot resulting a stronger chemical adsorption. When the electric field value further increased to 0.008 a.u., the distance between CO<sub>2</sub> and BC<sub>3</sub> suddenly decreased to 1.75 Å, the angle of  $\angle$ O–C–O further decreased to 127.3°, and the amount of charge transfer increased to 0.9 e. Seen from the Fig. 1d, the charge density between CO<sub>2</sub> and BC<sub>3</sub> increased dramatically, resulting that the interaction between the CO<sub>2</sub> molecular and the BC<sub>3</sub> surface has been greatly improved. These evidences indicate that under an applied electric field of 0.008 a.u., CO<sub>2</sub> is chemically adsorbed on the surface of BC<sub>3</sub>. It is

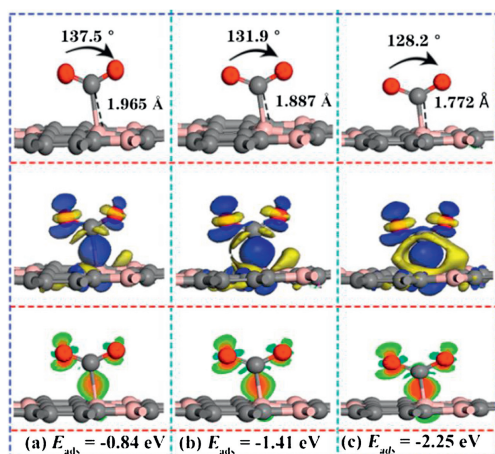


**Fig. 1.** Top and side views of the most favorable CO<sub>2</sub> adsorption configurations, corresponding charge density difference in an external E-field of (a) 0.002 a.u., (b) 0.004 a.u., (c) 0.006 a.u., (d) 0.008 a.u.. The blue (yellow) color represents accumulation (diminishing) of electrons. The values of isosurface of (a-c) are 0.003 e/Å<sup>3</sup>, and (d) is 0.02 e/Å<sup>3</sup>.

clearly that CO<sub>2</sub> is ‘fixed’ to BC<sub>3</sub> surface *via* employing an electric field. The results indicate the interaction between CO<sub>2</sub> and BC<sub>3</sub> surface is enhanced, which cause the stronger gas monitoring sensitivity.

In order to explain the turning point of CO<sub>2</sub> adsorption affected by the electric field, three values of electric field (0.0065, 0.007, 0.0075 a.u.) were performed to explore the corresponding adsorption structure, adsorption energy and the differential charge density (as shown in Fig. 2). As shown in Fig. S3a (Supporting information), the adsorption distance suddenly decreased from 1.97 Å under the electric field value of 0.0065 a.u., which was 33% shorter than that of values at 0.006 a.u. (2.94 Å). It is noted that the adsorption energy was increased to −0.84 eV, which is almost diploid compared with that of −0.43 eV. In addition, the angle of ∠O-C-O is reduced to 137.5°, due to the strong interaction of electrons accumulate between CO<sub>2</sub> and the substrate, which indicates that can be regarded as strong chemisorption. When the E-field value increased to 0.0075 a.u., the adsorption distance was further shortened to 1.772 Å, the angle of ∠O-C-O was further reduced to 128.2°, and the adsorption energy was increased to −2.25 eV, which indicates that the interaction between CO<sub>2</sub> and BC<sub>3</sub> substrate is further enhanced. The applied E-field will promote the polarization of the charged particles [56,57], which will lead to the enhanced adsorption of CO<sub>2</sub> on the surface of BC<sub>3</sub> [55]. As shown in Figs. S3a–c (Supporting information), the critical point of CO<sub>2</sub> from physical adsorption to chemical adsorption on the BC<sub>3</sub> substrate is between 0.006–0.0065 a.u.. As shown in Fig. S4, the minimum E-field value (0.0065 a.u.) required for the conversion of CO<sub>2</sub> from physical adsorption to chemical adsorption is much smaller than that of Penta-graphene [4], and is also smaller than the E-field value required in the previous literature to capture CO<sub>2</sub> materials, such as C<sub>2</sub>N [55], h-BN [27], C<sub>3</sub>N [58], P-doped C<sub>60</sub> [59], P-doped graphene [60]. It is a pity that the value is just slightly larger than that of MoS<sub>2</sub> [29] and PC<sub>5</sub> [24]. Such as it is, BC<sub>3</sub> captures CO<sub>2</sub> is more economical and efficient.

In order to further study the interaction between CO<sub>2</sub> and BC<sub>3</sub> substrate under different E-field values, partial density of states (PDOS) was studied and presented in Fig. S5 (Supporting information). As shown in Fig. S5, between the E-field value of 0 ~ 0.004 a.u., there is almost no overlap [48] between the orbits of CO<sub>2</sub> and B atoms, which indicates that the interaction between them is relatively weak physical adsorption. At this time, the E-field has little effect on the capture of CO<sub>2</sub>. However, when the

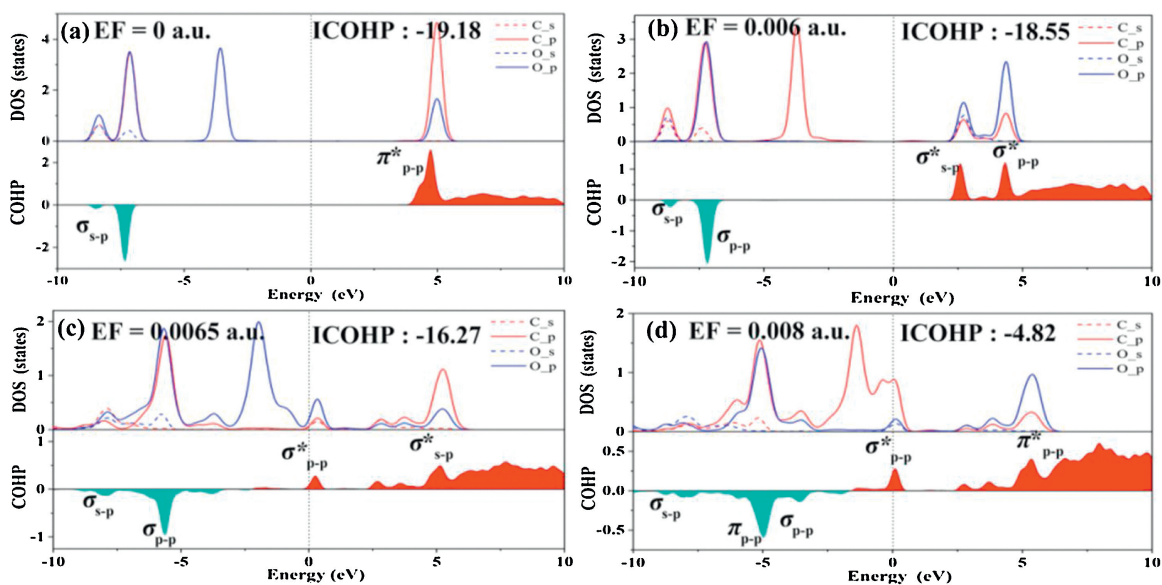


**Fig. 2.** The side view of the most stable configurations and corresponding charge density difference in E-field of (a) 0.0065 a.u., (b) 0.007 a.u., (c) 0.0075 a.u.. The blue (yellow) color represents accumulation (diminishing) of electrons. Red and green represent increasing and decreasing electron density. The values of isosurface of (a), (b), (c) are 0.015 e/Å<sup>3</sup>, respectively.

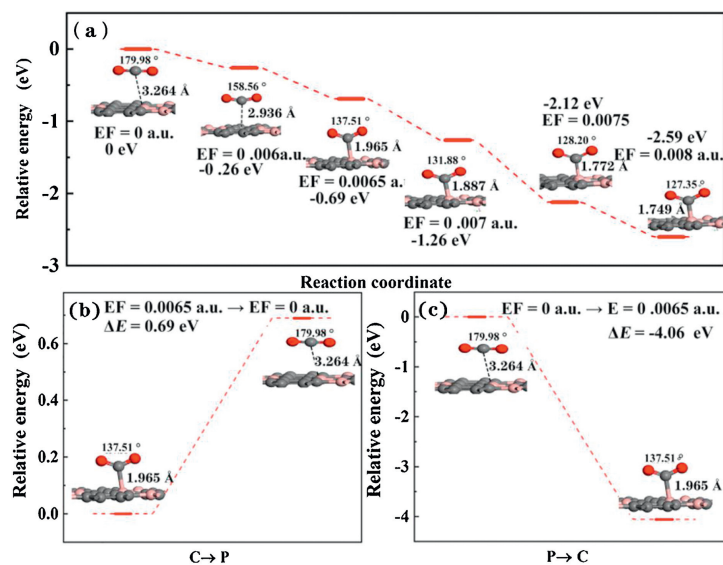
E-field value increases to 0.006 a.u., it can be seen that the overall PDOS has a significant right shift, and there is a clear overlapped orbitals of 2p (CO<sub>2</sub>) and 2p (B) between 3.0 ~ 5.0 eV, indicating the adsorption effect was enhanced, similarly as the previous results [52]. Similarly, when the E-field values at 0.0065 and 0.007 a.u., the orbital overlap formed by the 2p (CO<sub>2</sub>) and 2p (B) hybrid contributions can appear at −5.5, −1.2 and 5.0 eV, and the same results can be seen at 0.0075 and 0.008 a.u. It should be noted that the applied E-field can significantly enhance the ability of BC<sub>3</sub> to capture CO<sub>2</sub>.

In order to further study the mechanism of CO<sub>2</sub> adsorption on BC<sub>3</sub> under an external E-field, the crystal orbital hamiltonian population (COHP) was performed to analyze the different adsorption structures [61–64]. As shown in Fig. 3, the red area is the contribution of the anti-bonding state, and the cyan area is the bonding state. Fig. 3a shows that in the absence of an external E-field, the molecular orbital between C–O is composed by the  $\sigma_{s-p}$  bonding orbital (at −6.8 eV) that formed by the interaction of C<sub>s</sub> and O<sub>p</sub> and the anti-bonding  $\pi^*_{p-p}$  (at 4.5 eV) formed by the interaction of C<sub>s</sub> and O<sub>p</sub>. In addition, integrated COHP (ICOHP) was calculated to analyze the energy integral up to the highest occupied bands, and the value of ICOHP quantitatively reflected the strength of bonding. As shown in Figs. 3a–d, as the value of the E-field increases, ICOHP gradually decreases, which indicating the weaker of C–O bond in CO<sub>2</sub>. The reason is that with the increases of the E-field, the interaction between the C of CO<sub>2</sub> and B atom is enhanced, inducing the electron cloud more shared with the B atom and a reduction in the overlap of the electron cloud of the C–O, which is accordance with the weakened C–O interaction. Figs. S6a–d show that as the E-field strength increases, the ICOHP between C and B atom increases with the bonding state broaden, indicating that the interaction between C and B is strengthen. It should be highlighted that in the range of 0.0065–0.0075 a.u., the bonding molecular orbital  $\sigma_{p-p}$  between C and B atoms is mainly composed by C<sub>p</sub> and B<sub>p</sub>. Interestingly, bonding molecular orbitals  $\sigma_{p-p}$  which located in lower energy contributed by C<sub>p</sub> and B<sub>p</sub> appeared near the Fermi level.

In order to further explore the possible mechanism of CO<sub>2</sub> adsorption and desorption on BC<sub>3</sub> substrate under the external E-field existence, the potential energy [65] surface was shown in Fig. 4. The energy of the CO<sub>2</sub> adsorption system without an E-field is defined as 0. Clearly, the energy of the system gradually decreases along with E-field increases, indicating that the adsorption of CO<sub>2</sub> under the external E-field is a spontaneous exothermic process. In addition, when applying 0.006 a.u. E-field, the relative system energy decreases to −0.26 eV. It is worth noting that when the external E-field increases to 0.0065 a.u., the relative energy difference immediately decreases to −0.69 eV, well consistent with the  $E_{ads}$  of −0.84 eV, resulting to strong chemical adsorption. Besides, we also systematically studied the capture and separation of CO<sub>2</sub> gas on the BC<sub>3</sub> substrate by switching the E-field. As shown in Fig. 4b, under the external electric field of 0.0065 a.u., CO<sub>2</sub> is chemically adsorbed on the surface of BC<sub>3</sub>. When the external electric field is turned off, the CO<sub>2</sub> adsorption distance is extended from 1.965 Å to 3.264 Å with the CO<sub>2</sub> angle shifted from 137.5° to 179.9° after restructuring optimization. The above results demonstrate that the removal of the electric field can promote the transition from chemical adsorption to physical adsorption, in this process the endothermic reaction relative energy value is 0.69 eV. Similarly, when the E-field is switched on, CO<sub>2</sub> changed from physical adsorption to chemical adsorption with a linear structure to a V-shaped structure, and the adsorption distance is shortened to 1.965 Å. This process is exothermic process with relative energy −4.06 eV without energy barrier. These results shows that the adsorption/desorption of CO<sub>2</sub> on the BC<sub>3</sub> substrate can be easily controlled by the opening/closing of the electric field.



**Fig. 3.** PDOS (partial density of states) and molecular orbital of CO<sub>2</sub> molecule adsorbed on BC<sub>3</sub> under different E-field values, above and below the 0 value of crystal orbital hamiltonian population (COHP) are the anti-bonding state and bonding state, respectively. Red stands for bonding contributions, while cyan stands for antibonding contributions.



**Fig. 4.** The energy profile for the chemisorption to physisorption of the adsorbed CO<sub>2</sub> over BC<sub>3</sub> due to removing or switching of E-field (EF = 0.0065 a.u.).

CH<sub>4</sub> and H<sub>2</sub> are high-quality energy gases and important raw materials for the manufacture and synthesis of many chemical products. However, in the actual production and processing of CH<sub>4</sub> and H<sub>2</sub>, the presence of CO<sub>2</sub> impurities will inevitably be introduced. Therefore, the separation of CO<sub>2</sub> from the CO<sub>2</sub>/H<sub>2</sub>/CH<sub>4</sub> mixed gas is of great research significance for improving the purity of CH<sub>4</sub> and H<sub>2</sub>. In order to evaluate the separation effect of the electric field on the CO<sub>2</sub>/CH<sub>4</sub>/H<sub>2</sub> mixed gas, we studied the adsorption behavior of BC<sub>3</sub> to CO<sub>2</sub>/CH<sub>4</sub>/H<sub>2</sub> under the E-Field of 0.004 and 0.0065 a.u.. In the external E-Field of 0.004 a.u., the most stable adsorption structure is shown in Fig. S7 (Supporting information). The adsorption distances of CO<sub>2</sub>/H<sub>2</sub>/CH<sub>4</sub> on the BC<sub>3</sub> substrate are 3.032, 2.118, and 2.952 Å, and the corresponding  $E_{\text{ads}}$  are -0.35, -0.43 and -0.38 eV, respectively. No obvious electron distribution was observed between the CO<sub>2</sub>/CH<sub>4</sub> gas molecules and the substrate. According to the difference in adsorption energy, it

can be concluded that H<sub>2</sub> can be efficiently separated from the CO<sub>2</sub>/H<sub>2</sub>/CH<sub>4</sub> mixed gas under the 0.004 a.u.. In addition, the adsorption behaviors of CO<sub>2</sub>/H<sub>2</sub>/CH<sub>4</sub> gas under an electric field of 0.0065 a.u. has been studied, which is shown in Fig. S8 (Supporting information). Differ from the physical adsorption of CO<sub>2</sub>/H<sub>2</sub>/CH<sub>4</sub> under no electric field (Fig. S2.), the  $E_{\text{ads}}$  and  $D_{\text{m}}$  of CO<sub>2</sub> change significantly (-0.11 to -0.84 eV, 3.371 to 1.965 Å). There is no electron density distribution between the H<sub>2</sub>/CH<sub>4</sub> and the BC<sub>3</sub> substrate, but there is a distinct electron density overlaps between the CO<sub>2</sub> and the substrate, as shown in Fig. S8. Obviously, the adsorption properties of CO<sub>2</sub> and H<sub>2</sub>/CH<sub>4</sub> are different, under this circumstance the former is chemical adsorption and can be captured by BC<sub>3</sub>, but in fact the latter is physical adsorption.

We used the spin-polarized DFT calculation method to study the adsorption/desorption behavior of CO<sub>2</sub> gas with an external Z-axis vertical positive electric field. An external electric field

vertical to the BC<sub>3</sub> surface can significantly enhance the interaction between CO<sub>2</sub> and the substrate. In particular, in the range of 0.006–0.0065 a.u. electric field value, CO<sub>2</sub> shifted from BC<sub>3</sub> physisorption on C<sub>4</sub>B<sub>2</sub> vacancy with the rapid decrease of  $E_{\text{ads}}$  by 663% (–0.11 to –0.84 eV). Desorption of chemisorbed CO<sub>2</sub> occurred from the BC<sub>3</sub> surface without passing any energy barrier and released an energy of 0.69 eV when the E-Field is switched off. Utilizing the difference in the adsorption properties of CO<sub>2</sub>/CH<sub>4</sub>/H<sub>2</sub> on BC<sub>3</sub> under the E-Field value of 0.004 and 0.0065 a.u., H<sub>2</sub> and CO<sub>2</sub> gas can be separated from the CO<sub>2</sub>/H<sub>2</sub>/CH<sub>4</sub> gas mixture. These results confirmed that BC<sub>3</sub> is a potential material for capturing CO<sub>2</sub> under external electric field.

### Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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### Appendix A. Supplementary data

Supplementary material related to this article can be found, in the online version, at doi:<https://doi.org/10.1016/j.ccllet.2021.03.038>.

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