



Communication

Thermo-induced structural transformation with synergistic optical and magnetic changes in ytterbium and erbium complexes

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ABSTRACT

Dinuclear ytterbium and erbium based bifunction complexes $\text{Ln}_2\text{L}_2(\text{depma}_2)\text{Cl}_2$ (**1-Ln**, Ln = Yb and Er, $\text{H}_2\text{L} = \text{N}^1, \text{N}^3$ -bis(salicylideneimino)diethylenetriamine, depma_2 = dimerized 9-diethyl-phosphonomethylanthracene) are reported. They undergo thermo-induced consecutive phase transitions, first the dissociation of depma_2 ligand forming $\text{LnL}(\text{depma})\text{Cl}$ (**2-Ln**) and then the release of chloroethane forming $\text{LnL}(\text{epma})$ (**3-Ln**, $\text{epma} = 9$ -ethylphosphonomethylanthracene). The structural transformations are accompanied with synergetic switch of the luminescence in visible and NIR regions and also magnetic dynamics.

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Combining two properties into one molecule is an effective way to obtain bifunctional molecular materials. Lanthanide compounds are excellent candidates for the construction of luminescent single-molecule magnets (SMMs) because they exhibit characteristic luminescence with well-resolved emission bands and long lifetimes and, meanwhile, show large single-ion magnetic anisotropy due to the presence of large unquenched orbital angular momentum and strong spin-orbit coupling [1–3]. While a number of lanthanide complexes were synthesized showing coexistence of luminescence and SMM behavior, manipulation of the two functions in a synergistic manner remains to be a challenging task. Only a few dysprosium complexes were reported exhibiting simultaneous switch of luminescence and magnetic dynamics through mechanical force [4,5], light irradiation [6–8] and thermal treatment [9–11]. The erbium and ytterbium complexes are similar to the dysprosium ones in some way because they also possess Kramers ground states with strong magnetic anisotropy [12–16]. They further distinct themselves from the dysprosium ones in their characteristic emission in near infrared (NIR) instead of visible region which feature is appealing for applications in material science and biotechnologies [17]. However, although luminescent Er- and Yb-SMMs are not uncommon [1–3,18–22], synergistic switching of magnetic and

optical properties has never been documented for Er- and Yb-based compounds.

Herein we report two isostructural dinuclear complexes $\text{Ln}_2\text{L}_2(\text{depma}_2)\text{Cl}_2$ (**1-Ln**, Ln = Yb and Er) incorporating a thermo-responsive pre-photodimerized 9-diethyl-phosphonomethylanthracene (depma_2) as bridging ligand, and N^1, N^3 -bis(salicylideneimino)diethylenetriamine (H_2L) as terminal ligand, as an extension of our previous work on dysprosium complex [11]. Upon heating, the depma_2 ligand in **1-Ln** undergoes thermo-induced dedimerization at ca. 170 °C forming mononuclear species $\text{LnL}(\text{depma})\text{Cl}$ (**2-Ln**), followed by thermo-induced elimination of $\text{CH}_3\text{CH}_2\text{Cl}$ at about 250 °C forming $\text{LnL}(\text{epma})$ (**3-Ln**, $\text{epma} = 9$ -ethylphosphonomethylanthracene). The two consecutive structural transformations are accompanied with the turn-on of photoluminescence in visible and/or NIR regions originating from the anthracene excimer and lanthanide ion, respectively, and acceleration of magnetic relaxation at low temperature. To our knowledge, **1-Ln** are the first examples of Yb- and Er-based compounds showing synergetic switching of magnetic and optical properties upon external stimulus.

Complexes **1-Ln** were synthesized by reacting $\text{LnCl}_3 \cdot 6\text{H}_2\text{O}$ with depma_2 , H_2L and Et_3N in methanol solution at room temperature. The two complexes are isostructural and crystallize in the orthorhombic space group $Pbca$ (Table S1 in Supporting information). **1-Yb** is selected as a representative for a detailed structural description. The asymmetric unit of **1-Yb** consists of one Yb^{III} , one L^{2-} , one Cl^- and half depma_2 . The Yb^{III} ion is ligated by two oxygen (O4, O5) and three nitrogen (N1, N2, N3) atoms from L^{2-} in the

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equatorial plane [Yb1-O: 2.156(3)–2.163(3) Å, Yb1-N: 2.486(3)–2.517(3) Å], and one oxygen (O1) from depma₂ and one Cl[−] (Cl1) anion in the axial positions [Yb1-O1: 2.233(2) Å, Yb1-Cl: 2.604(1) Å] (Fig. 1, Fig. S3 and Table S2 in Supporting information). The geometry of the {YbN₃O₃Cl} core is best described as a pentagonal bipyramid, according to the Continuous Shape Measure (CShM = 0.498) analyses (Table S3 in Supporting information) [23]. The metal centers are related through a center of crystallographic inversion symmetry to form a dinuclear molecule (Fig. 1a). The Yb···Yb distance within the dimer is 15.193(1) Å. Neighboring dimers are linked by N—H···Cl hydrogen bonds (N···Cl: 3.399 Å) to form a supramolecular layer (Fig. 1b and Table S4 in Supporting information). These layers are further stacked in staggered manner along the *b*-axis and stabilized by van der Waals interactions (Fig. S3). The shortest intermolecular Yb···Yb distance is 7.003(5) Å.

The coordination geometry of Er in **1-Er** is identical to that of Yb. The geometry of the {ErN₃O₃Cl} core is also best described as a pentagonal bipyramid (CShM = 0.487) (Table S3). The bond distances are Er1-O: 2.160(3)–2.236(3) Å, Er1-N: 2.486(3)–2.517(4) Å and Er1-Cl: 2.605(1) Å, respectively. The shortest intermolecular Er···Er distance is 7.030(7) Å (Fig. S4 and Table S2 in Supporting information).

As the pre-photodimerized dianthracene derivatives can cleave to monomers upon heating, thermal analyses were conducted for **1-Ln**. Although the weight keeps almost constant in the TG curves below 200 °C (Fig. 2a), the DSC curves show an exothermic peak at 168 °C for **1-Yb** ($\Delta H = 47.79$ kJ/mol) and **1-Er** ($\Delta H = 44.27$ kJ/mol) (Fig. 2b). This is ascribed to the dissociation of depma₂ in **1-Ln**, proved by *in-situ* variable temperature infrared (IR) spectra. As shown in Figs. S5 and S6 (Supporting information), upon heating **1-Ln** to 175 °C, peaks at 1250 and 878 cm^{−1} emerge corresponding to the C—H in-plane and out-of-plane bending vibrations of anthracene

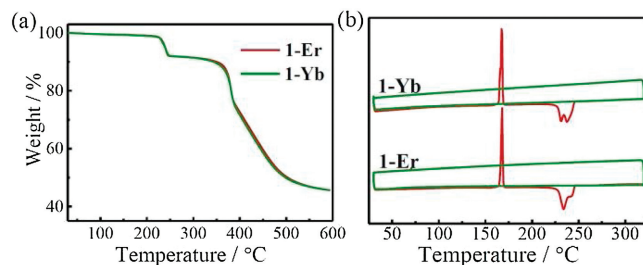


Fig. 2. (a) TG and (b) DSC curves for complexes **1-Yb** and **1-Er**, measured under nitrogen atmosphere at a heating rate of 5 °C/min.

[24]. The dissociation of depma₂ in **1-Ln** leads to the formation of monomeric species LnL(depma)Cl (**2-Ln**). Moreover, the DSC curves show another endothermic doublet peak at 231 and 237 °C for **1-Yb** and 234 and 240 °C for **1-Er**, with the total enthalpy values of 35.3 and 32.8 kJ/mol, respectively. TG analyses suggest the release of one molecule of CH₃CH₂Cl in the range of 200–250 °C for both **1-Yb** and **1-Er** (obs. 7.4%, calcd. 7.8%) (Fig. 2a), which is further confirmed by the TG-MS analysis (Fig. S7 in Supporting information). The release of CH₃CH₂Cl is a result of intramolecular solid-state reaction between the ethyl group of depma and the coordinated Cl[−] anion at elevated temperature, which phenomenon is unusual and was previously observed only in rare cases of coordination systems [11,25]. The elimination of chloroethane in **2-Ln** is associated with the broken of one C—O(P) bond of depma, leading to LnL(epma) (**3-Ln**, epma[−] = 9-ethylphosphonomethylanthracene anion). This process can be monitored by *in-situ* IR spectra of **1-Ln**. Indeed, the peak intensity of the P—O stretching vibration at 1016 cm^{−1} decreases with increasing temperature, and finally the peak disappears above 200 °C (Figs. S5 and S6). The molecular formulae of **2-Ln** and **3-Ln** were confirmed by elemental analyses. The powder XRD measurements corroborate that the structures of **1-Ln**, **2-Ln** and **3-Ln** are different from each other (Fig. S8 in Supporting information), but identical to the corresponding Dy and Gd analogues [11]. The single crystallinity of **1-Ln** is not retained after thermal treatment.

The diffuse reflectance spectra of all of the complexes show absorption bands in the range of 200–450 nm, originating from the π – π^* transitions of the ligands (Figs. S9 and S10 in Supporting information). Besides, **1-Yb**, **2-Yb** and **3-Yb** show the Yb³⁺ f-f absorption peaks appearing at 912 and 968 nm, ascribed to the transition of $^2F_{7/2}$ – $^2F_{5/2}$ [24]. For **1-Yb** and **2-Yb**, the peak at 912 nm is more intensive, while for **3-Yb**, the peak at 968 nm is more intensive (Fig. S9). For **1-Er** and **2-Er**, similar Er³⁺ f-f absorption peaks appear at 453 ($^4I_{15/2}$ – $^4F_{5/2}$), 490 ($^4I_{15/2}$ – $^4F_{7/2}$), 521 ($^4I_{15/2}$ – $^2H_{11/2}$), 530, 544, 548 ($^4I_{15/2}$ – $^4S_{3/2}$), 646, 656, 664, 675 ($^4I_{15/2}$ – $^4F_{9/2}$) and 796 nm ($^4I_{15/2}$ – $^4I_{9/2}$) [26]. Notably, **3-Er** shows the same transition but the absorption intensities are slightly changed (Fig. S10). The slight differences in absorption bands of **3-Ln** from those of **1-Ln** and **2-Ln** are related to the change in the coordination environment of Ln in **3-Ln** due to the release of coordination Cl[−].

Luminescence spectra of all complexes were measured in order to study the effect of structural changes on the photophysical properties. For Yb-complexes, when excited at 310 nm, **1-Yb** is not emissive in the visible region, but emissive in the NIR region (Fig. 3a). It exhibits well-resolved multiline emission bands at 979 and 1041 nm, in which the former is the most intensive one. The emission bands are attributed to the transition from the first excited state $^2F_{5/2}$ to the manifold $\pm m_j$ sublevels of the ground state $^2F_{7/2}$ of the Yb^{III} ion as a consequence of ligand-field effect [27]. The quench of the ligand emission together with the observation of the characteristic NIR emission of Yb^{III} ion suggest an efficient energy transfer from the ligand to Yb^{III} ion. For **2-Yb**, characteristic NIR emission is again observed. Unlike **1-Yb**,

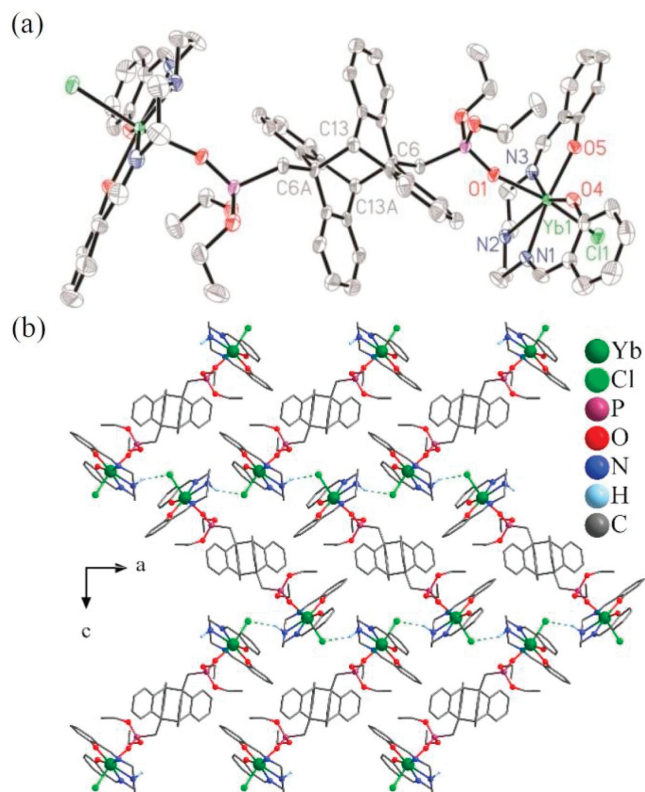


Fig. 1. (a) Molecular structure of complex **1-Yb** with an atomic labelling scheme (50% probability). (b) Supramolecular layer structure of **1-Yb** formed by hydrogen bonds.

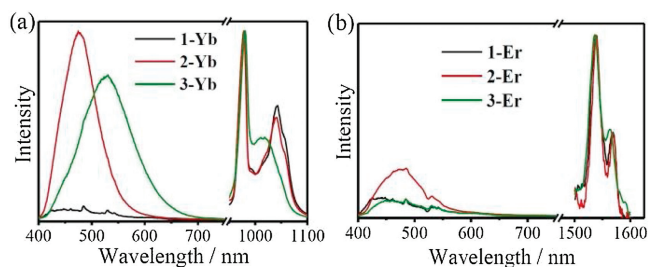


Fig. 3. The solid-state emission spectra (excited at 310 nm) for complexes **1-Ln**, **2-Ln** and **3-Ln** at visible and NIR regions.

however, it also shows a broad emission at 476 nm, attributed to the formation of excimer of staggered π - π stacked anthracene groups in the structure [28]. The quantum yield (QY) and average lifetime (τ_{av}) are 1.57% and 4.34 ns, respectively (Fig. S13a and Table S5 in Supporting information). Complex **3-Yb** also exhibits dual emissions in the visible and NIR regions. In this case, the ligand-centered emission peak is red-shifted to 530 nm (QY = 1.40%, τ_{av} = 10.84 ns) (Fig. S13b in Supporting information), corresponding to the excimer of face-to-face π - π stacked anthracene groups [6,11]. In the NIR region, **3-Yb** shows the emission at 981 and 1011 nm, also ascribed to the ${}^2F_{5/2} \rightarrow {}^2F_{7/2}$ transition of Yb^{III}. The results indicate that energy transfer occurs from the triplet state of ligand to both the Yb^{III} ion and the anthracene excimer in **2-Yb** and **3-Yb**. The ${}^2F_{5/2}$ (Yb) lifetimes are explored for **1-Yb**, **2-Yb** and **3-Yb**. The emission decay profiles can be well fitted with single exponential functions, yielding lifetime of 4.2–4.4 μ s (Fig. S16 and Table S6 in Supporting information).

For Er-complexes, although weak emission bands are observed in visible region for all three complexes, their quantum yields are close to zero. However, they exhibit emission in NIR region at 1538 and 1564 nm (Fig. 3b and Fig. S12 in Supporting information), corresponding to the ${}^4I_{13/2} \rightarrow {}^4I_{15/2}$ transition of Er^{III} ion [27]. Noting that emission in visible region is found for **2-Yb** and **3-Yb** but not for **2-Er** and **3-Er**, an efficient energy transfer from the π - π stacked anthracene excimer to Er^{III} ion could occur. This is in agreement with the fact that the energy gap between the ground and excited state of Er^{III} ion (${}^4I_{13/2}$ to ${}^4I_{15/2}$) is lower than that of Yb^{III}

ion (${}^2F_{5/2}$ to ${}^2F_{7/2}$). However, the ${}^4I_{13/2}$ (Er) lifetime is not detectable for all Er complexes (Fig. S17 in Supporting information).

The structural transformation of **1-Ln** upon heating is also reflected by their magnetic properties. The temperature dependent dc magnetic susceptibilities (Fig. S18 in Supporting information) reveal that the room temperature $\chi_M T$ values (per Ln) of 2.50 (**1-Yb**), 2.50 (**2-Yb**), 2.49 (**2-Yb**), 11.18 (**1-Er**), 11.18 (**2-Er**), and 11.19 (**3-Er**) $\text{cm}^3 \text{K mol}^{-1}$ agree well with the theoretical values of 2.57 $\text{cm}^3 \text{K mol}^{-1}$ for Yb^{III} and 11.48 $\text{cm}^3 \text{K mol}^{-1}$ for Er^{III}. In all cases, the $\chi_M T$ decreases progressively upon cooling, attributed to the depopulation of the M_J levels and possible weak antiferromagnetic interactions. The unsaturation of the magnetization as well as the non-superimposition of M vs. H/T plots suggest the presence of significant magnetic anisotropy and/or low-lying excited states in **1-Ln**, **2-Ln** and **3-Ln** (Figs. S19–S24 in Supporting information).

To investigate the magnetic dynamics of these complexes, ac magnetic susceptibilities were performed at low temperatures. All show field-induced magnetic relaxations except **3-Er** which does not show any in-phase (χ') and out-of-phase (χ'') signals even under an external field of 2 kOe (Fig. 4 and Figs. S25–S30 in Supporting information). For Yb-complexes, the χ' and χ'' vs. frequency plots were measured under a dc field of 1 kOe. The relaxation time (τ) was extracted by fitting the Cole-Cole plots using a generalized Debye model (Tables S7–S9 in Supporting information). The $\ln \tau$ vs. T^{-1} curve can be best fit by considering the direct and Orbach processes using Eq. 1 (Fig. 4d) [29], where U_{eff} represents the energy barrier. The resulted parameters are $U_{eff} = 23.5 \text{ K}$, $\tau_0 = 1.2 \times 10^{-6} \text{ s}$ and $A = 13.5 \text{ s}^{-1} \text{ K}^{-1}$ for **1-Yb**, $U_{eff} = 22.3 \text{ K}$, $\tau_0 = 9.1 \times 10^{-7} \text{ s}$ and $A = 161.9 \text{ s}^{-1} \text{ K}^{-1}$ for **2-Yb** and $U_{eff} = 13.1 \text{ K}$, $\tau_0 = 9.4 \times 10^{-7} \text{ s}$ and $A = 1787.2 \text{ s}^{-1} \text{ K}^{-1}$ for **3-Yb**. Compared with **1-Yb**, the energy barrier of **2-Yb** is quite close indicating that the dissociation of dianthracene poses little influence on the coordination geometry of Yb^{III} ion. However, the energy barrier of **3-Yb** drops nearly a half, which is consistent with the change of coordination geometry of Yb^{III} ion due to the release of coordination Cl⁻ anion. These U_{eff} values are comparable to those of the other Yb^{III}-based complexes showing field-induced SMM behavior (2.0–54.7 K) [1,2,13,18–22,30].

$$\tau^{-1} = AT + \tau_0^{-1} \exp\left(-\frac{U_{eff}}{kT}\right) \quad (1)$$

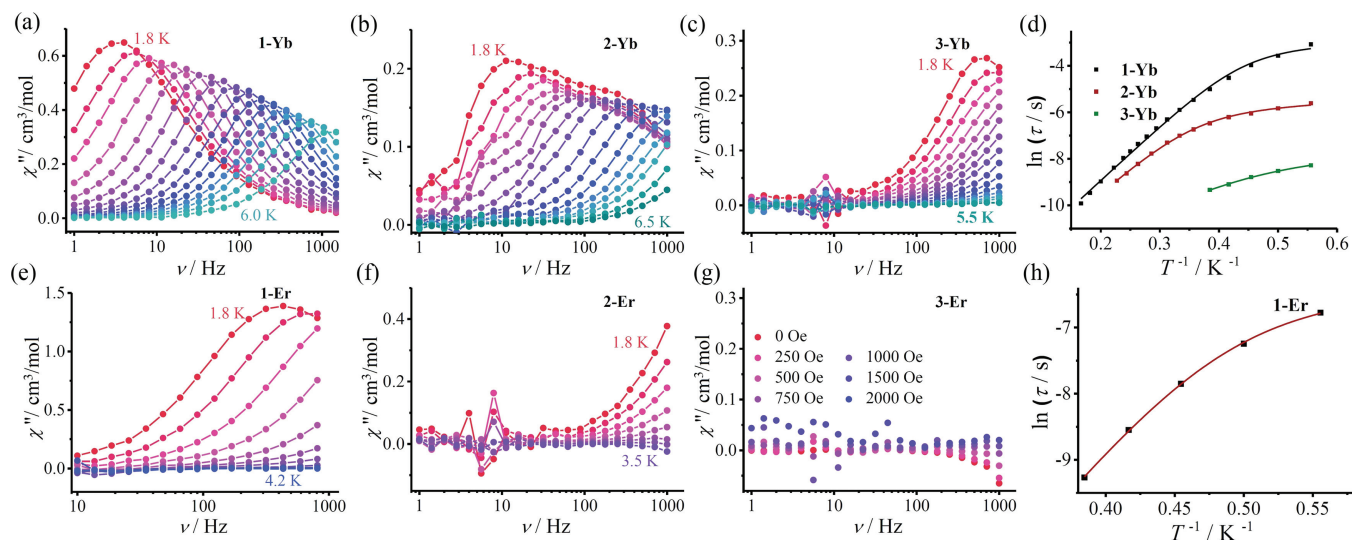


Fig. 4. Frequency dependence of χ'' signals measured at 1 kOe dc field for (a) **1-Yb**, (b) **2-Yb** and (c) **3-Yb** at the depicted temperatures. (d) Plots of $\ln \tau$ vs. T^{-1} for **1-Yb**, **2-Yb** and **3-Yb**. Frequency dependence of χ'' signals measured at 0.75 kOe dc field for (e) **1-Er** and (f) **2-Er** at the depicted temperatures. (g) Frequency dependence of χ'' signals of **3-Er** in the indicated dc fields at 2.0 K. (h) Plots of $\ln \tau$ vs. T^{-1} for **1-Er**.

For **1-Er**, clear frequency dependent peaks for χ'' were observed under a dc field of 0.75 kOe (Fig. 4e and Fig. S28). The $\ln\tau$ vs. T^{-1} curve can also be fitted by Eq. 1 (Fig. 4h), with the parameters of $U_{\text{eff}} = 24.6$ K, $\tau_0 = 7.9 \times 10^{-9}$ s and $A = 415.3$ s $^{-1}$ K $^{-1}$. The energy barrier is within the range of those previously reported for field-induced Er^{III} SMM complexes (7.5–34.8 K) [1,2,18,20,30]. For **2-Er**, the out of phase ac susceptibility data exhibit the frequency dependence below 3.5 K, but no maximum appears (Fig. 4f and Fig. S29). Therefore, we cannot obtain the energy barrier of **2-Er** from the Eq. 1. Moreover, complex **3-Er** does not show any ac signals (Fig. 4g and Fig. S30). The significant decrease of energy barrier of **2-Er** comparing to that of **1-Er** and the switch-off of slow magnetic relaxation behavior of **3-Er** suggest that the structural transformations associated with the dissociation of dianthracene moieties and the release of CH₃CH₂Cl could have more significant influence on the magnetic anisotropy of Er^{III} than that of Yb^{III} ion.

In summary, we report the first examples of Yb- and Er-complexes incorporating a pre-photodimerized dianthracene ligand, Ln₂L₂(depma₂)Cl₂ (**1-Ln**, Ln = Yb and Er). Both experience a two-step thermo-induced structural transformation, first the dissociation of dianthracene and then the release of CH₃CH₂Cl. The consecutive phase transitions are accompanied by simultaneous change of the photoluminescence in the visible and NIR regions at room temperature and also magnetic dynamics at low temperature. This work provides a feasible route to the design of bifunctional Yb- and Er-based materials with synergetic and switchable magneto-optical properties. These properties together with their ability in emitting NIR luminescence may have potential applications in material science and biotechnology.

Declaration of competing interest

The authors report no declarations of interest.

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Appendix A. Supplementary data

Supplementary material related to this article can be found, in the online version, at doi:<https://doi.org/10.1016/j.ccl.2020.10.019>.

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